

16th US-Japan Workshop on Fusion Neutron Sources for Nuclear Assay and Alternate Applications

Madison, WI September 30th – October 2nd, 2014



Simulations for Multiple-grid IEC

Drew Chap

NASA Space Technology Research Fellow

Graduate Student, Aerospace Engineering, University of Maryland, College Park



Multiple-grid IEC – brief history

Sedwick et al. used additional grids to focus ion beams and increase ion confinement time

3 TAKEAWAYS:

- **1: Ion lifetimes extended:** From 10's of passes to 10³-10⁶ passes
- Greater confinement time
 + Counter-stream instability
 + IEC trap kinematics
 = Ion bunching





Multiple-grid IEC – current research

DODECAHEDRAL GRIDS

- 12 Faces → 6 beamlines
- Highly symmetric
- Another possibility: **Truncated Icosahedron** (Soccer Ball)
- Feed-throughs?

ION BUNCHING

• Potential well can be shaped to encourage ion bunch cohesion

MAGNETIC CORE

• Confinement of electrons in the core









2-GRID IEC





4-GRID IEC



Particle-particle Discrete Event Simulation

- Inter-particle forces are calculated directly (N-body simulation)
 - No need to solve Poisson's equation at each time step
- No global time-step, each particle is assigned its own time-step depending on its velocity and acceleration
 - Coulomb collisions are modeled directly by decreasing the time-step values of colliding particles.
- Static E&M fields are calculated once at the beginning of the simulation

$$\vec{a}_i = -\frac{1}{4\pi\epsilon_0} \frac{q_i}{m_i} \sum_{j \neq i} \frac{q_j}{r_{ij}^3} \vec{r}_{ij} + \frac{q_i}{m_i} \vec{E}$$



2-GRID Particle-particle simulation LOW DENSITY BUNCH









2-GRID Particle-particle simulation

What happens if we increase the density of the ion bunch?

2-GRID Particle-particle simulation

HIGHER DENSITY BUNCH



4-GRID Particle-particle simulation SAN

SAME DENSITY AS PREVIOUS SLIDE



Ion Bunching – The Kinematic Criterion



Ion Bunching – The Kinematic Criterion

Kinematic criterion:
$$\frac{dT}{dE} \ge 0$$

But $\frac{dT}{dE}$ can't be too large either!

Conditions have to be just right for ions to coalesce into bunches

FUTURE WORK:

"Sculpting" the IEC well to encourage bunch cohesion



Electron confinement in the IEC core









Approximate E&M fields along a beampath



Confinement of a single electron

Confinement of many electrons









Disadvantages of the particle-particle discrete event simulation

- Computation time scales as N²
- Only suitable (at this point) for modeling one species at a time (ions or electrons) for short timescales
- To model both species at once we need a hybrid PIC model

Hybrid Particle-in-cell model

 $\frac{\text{IONS} \rightarrow \text{PARTICLES}}{\text{ELECTRONS} \rightarrow \text{FLUID}}$

Long-timescale simulation requires time-steps based on the ion motion



Assume electrons reach a **thermalized steady-state** at each time-step $\left(\frac{dn_e}{dt} = 0\right)$ Hybrid Particle-in-cell model

Currently using the following governing equations:

Interested in steady-state solution

 ∇

CONTINUITY

$$\frac{\partial n_e}{\partial t} + \nabla \cdot (n_e \vec{v}) = S$$

VELOCITY

$$\vec{v} + (\vec{v} \cdot \nabla)\vec{v} = \frac{e}{m_e}(\nabla \Phi + \vec{v} \times \vec{B})$$

POISSON'S

$$e^{2}\Phi = \frac{e}{\epsilon_{0}}(n_{e} - n_{i})$$

But first we'll test the time-stepping solution

19

-3

Electron fluid model – TEST PROBLEM (not IEC) x 10 Wire locations **Electron source** 16 0.9 (perpendicular to plain) $(\#/m^{3}/s)$ 0.8 0.8 14 \otimes 0.7 12 0 0.6 0.6 10 Ο 0.5 8 Ø 0.4 0.4 6 \otimes 0.3 Θ 4 0.2 0.2 2 0.1 0 0 0 0.2 0.4 0.6 0.8 0 x 10 1 **B-field direction B-field magnitude (T)** 2 0.9 0.8 0.8 0.7 1.5 0.6 0.6 0.5 1 0.4 0.4 0.3 0.5 0.2 0.2

01

Simple 2D test problem:

- Magnetic field created by current-carrying wires
- Electron Source in Center



Simple 2D test problem:

- Magnetic field created by current-carrying wires
- Electron Source in Center

Comparison with particle-in-cell model with same conditions



Simple 2D test problem:

- Magnetic field created by current-carrying wires
- Electron Source in Center

Comparison with particle-in-cell model with same conditions



With increased B-field











The End



DENSITY
$$\frac{\partial n}{\partial t} + \frac{\partial}{\partial x}(nv) = S$$
VELOCITY
$$\frac{\partial v}{\partial t} + \frac{\partial}{\partial x}(v^2) = \frac{e}{m}\frac{\partial}{\partial x}\Phi$$
POTENTIAL
$$\frac{\partial^2}{\partial x^2}\Phi = \frac{e}{\epsilon_0}n$$

Using the **Roe scheme** (basically a selective upwind scheme) for density and velocity:

$$\frac{\partial u_i}{\partial t} + \frac{1}{2\Delta x} \left\{ (\Gamma_{i+1} - \Gamma_{i-1}) - \left| \frac{\Gamma_{i+1} - \Gamma_i}{u_{i+1} - u_i} \right| (u_{i+1} - u_i) + \left| \frac{\Gamma_i - \Gamma_{i-1}}{u_i - u_{i-1}} \right| (u_i - u_{i-1}) \right\} = S_i$$

Central differencing

Upwind correcting terms

For STEADY STATE: can't take the derivatives of absolute values:

$$\frac{\partial u_i}{\partial t} + \frac{1}{2\Delta x} \left\{ (\Gamma_{i+1} - \Gamma_{i-1}) - \left| \frac{\Gamma_{i+1} - \Gamma_i}{u_{i+1} - u_i} \right| (u_{i+1} - u_i) + \left| \frac{\Gamma_i - \Gamma_{i-1}}{u_i - u_{i-1}} \right| (u_i - u_{i-1}) \right\} = S_i$$

Use approximation: $|x| \approx \sqrt{x^2 + a^2}$ with small enough a

In our case, *a* is a velocity